

Quantum Raychaudhuri Equation

Saurya Das
University of Lethbridge
CANADA

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Plan

- Singularities and Raychaudhuri equation
- Quantum Raychaudhuri equation
- Implications for singularities
- Application: Cosmology

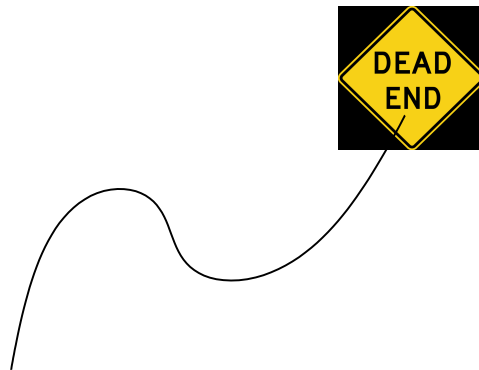
S. Das, Phys. Rev. **D89** 084068(2014) [arXiv:1311.6539]

A. F. Ali, S. Das, arXiv:1404.3093

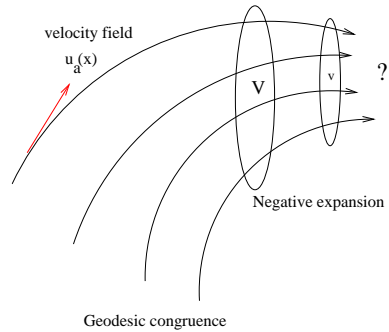
S. Das, arXiv:1405.4011

Singularities in spacetime

- Signal breakdown of General Relativity, start/end of time (big-bang/big-crunch), black holes, ...
- Curvature $\rightarrow \infty$ is indicative, but neither necessary, nor sufficient condition for singularities/pathological spacetime (coordinate dependent, many counterexamples)
- **Geodesic incompleteness** normally assumed to be necessary and sufficient condition for singular spacetimes



Raychaudhuri equation



Expansion $\theta = \frac{d \ln V}{d\lambda}$ (< 0 in above figure) ($\lambda =$ affine parameter)

$$\frac{du_{a;b}}{d\lambda} = u_{a;b;c} u^c = [u_{a;c;b} + R_{cba}{}^d u_a] u^c \quad (1)$$

$$= \underbrace{(u_{a;c} u^c)_{;b}} - u^c{}_{;b} u_{a;c} + R_{cba}{}^d u^c u_d \quad (2)$$

$= 0$ (geodesic equation)

$$= -u^c{}_{;b} u_{a;c} + R_{cbad} u^c u^d . \quad (3)$$

$$h_{ab} = g_{ab} - u_a u_b \text{ (induced 3-metric)}$$

$$u_{a;b} = \frac{1}{3}\theta h_{ab} + \sigma_{ab} + \omega_{ab} = \text{Trace} + \text{Traceless symmetric} + \text{anti-symmetric}$$

expansion shear twist

$$h^{ab} \frac{du_{a;b}}{d\lambda} = \text{Tr} \left(\frac{du_{a;b}}{d\lambda} \right) = \text{Tr} \left(-u^c{}_{;b} u_{a;c} + R_{cbad} u^c u^d \right)$$

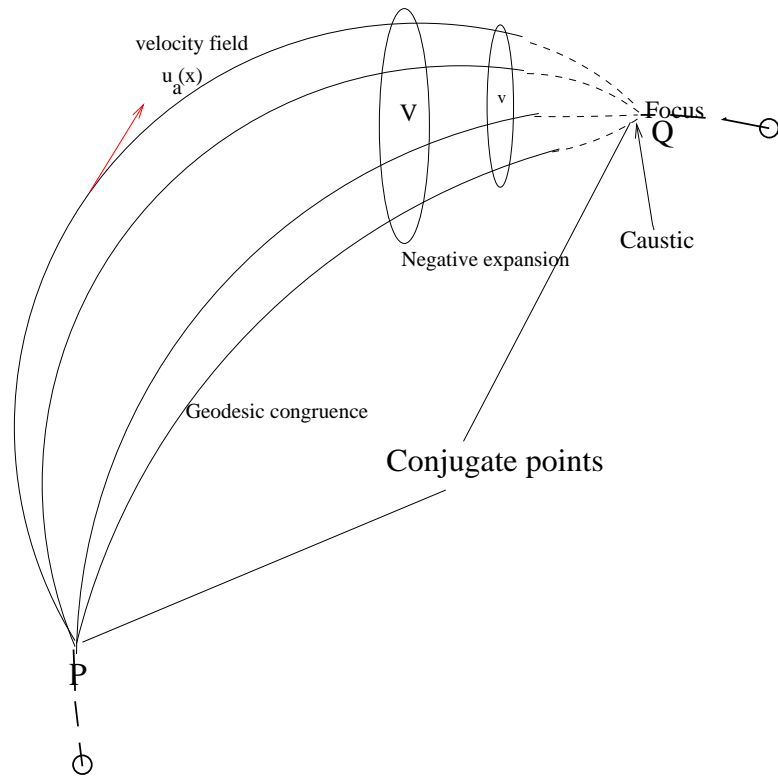
$$\frac{d\theta}{d\lambda} = -\frac{1}{3}\theta^2 - \sigma_{ab}\sigma^{ab} + \underbrace{\omega_{ab}\omega^{ab}}_{=0} - \underbrace{R_{cd}u^c u^d}_{>0} < 0$$

hypersurface orthog. strong energy cond.

If $\theta_0 = \theta(0) < 0$ (initially converging)

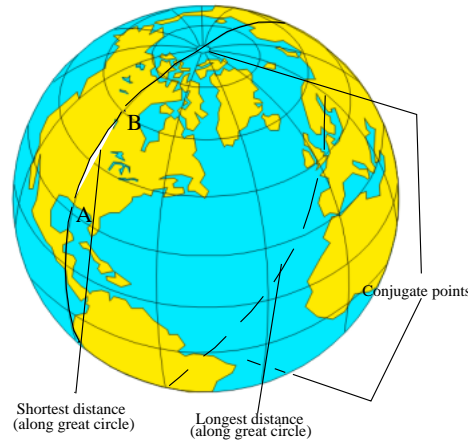
Focus/caustic for $\lambda \leq \frac{3}{|\theta_0|}$ (*finite proper time!*)

Raychaudhuri equation (1955) (also, Landau & Lifshitz)



Singularity theorems

- Conjugate points due to Raychaudhuri equation.
- Geodesics are no longer maximal length curves.



- Maximal geodesics predicted by global arguments, on the other hand.
- Sufficiently long geodesics cannot exist. Geodesics are incomplete.
- Singularities! (Singularity theorems) - for most spacetimes
- R. Penrose, Phys. Rev. Lett. **14** (1965) 57-59; S. W. Hawking, R. Penrose, Proc. Roy. Soc. Lond. **A314** (1970) 529-548.

Note

- *Generality*: for all reasonable spacetimes (gravity universal & attractive)
- Fluid picture: velocity field $u^a(x)$
- But, classical \rightarrow quantum? (Expectation values, Ehrenfest type theorem?)
- First, find a ‘quantum velocity field’

Non-relativistic limit

$$u^a(x) \rightarrow v^a(\vec{x}, t) \quad (a = 1, 2, 3), \quad u^0 = 1, \quad \lambda \rightarrow t$$

$$R_{cd}u^c u^d \rightarrow \nabla^2 V = 4\pi G\rho \geq 0,$$

$$\frac{d\vec{v}}{dt} = -\vec{\nabla}V(\vec{x}, t)$$

$$\frac{d\theta}{dt} = -\frac{1}{3}\theta^2 - \sigma_{ab}\sigma^{ab} - \nabla^2 V < 0 \quad \leftarrow \text{(focusing in finite time)}$$

Quantum fluid picture

$$\psi(\vec{x}, t) = \mathcal{R}e^{iS}$$

(Normalizable, single-valued, $\mathcal{R}, S = \text{Real}$.)

E.g. complete set of H-atom bound states and scattering wavefunctions, $e^2/4\pi\epsilon_0 \rightarrow GMm$)

$$\vec{v}(\vec{x}, t) = \frac{d\vec{x}}{dt} \equiv \frac{\hbar}{m} \mathcal{I}m \left(\frac{\vec{\nabla}\psi}{\psi} \right) = \frac{\hbar}{m} \vec{\nabla}S(\vec{x}, t) \leftarrow \text{quantum velocity field!}$$

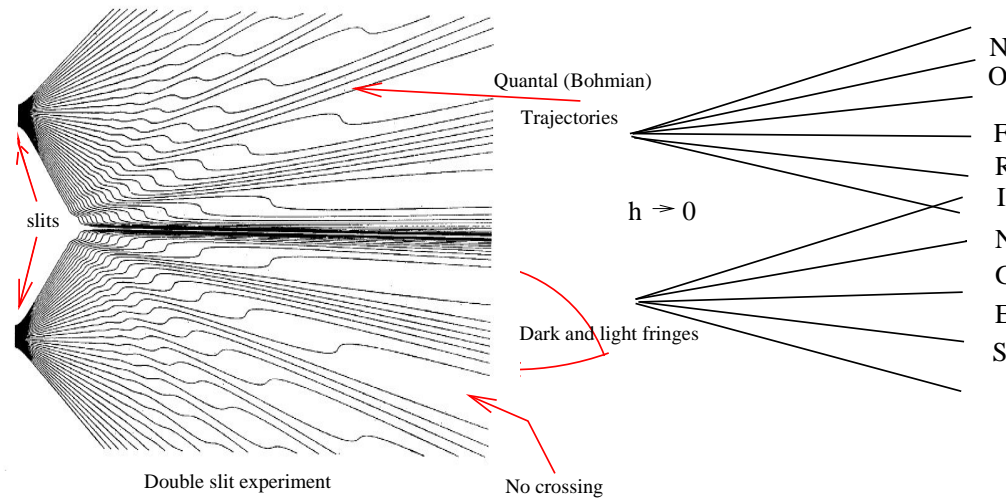
$$-\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi = i\hbar \frac{\partial \psi}{\partial t}$$

Real and imaginary parts \rightarrow

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0 \quad (\text{Probability conservation})$$

$$m \frac{d\vec{v}}{dt} = -m \vec{\nabla}V + \underbrace{\frac{\hbar^2}{2m} \vec{\nabla} \left(\frac{1}{\mathcal{R}} \nabla^2 \mathcal{R} \right)}_{-\vec{\nabla}V_Q} \quad (\text{Newton's law} + \text{quantum potential } V_Q!)$$

- Initially, particles distributed as $\rho(0) = |\psi(0)|^2$ ('quantum equilibrium')
- Prob. conservation \Rightarrow they remain distributed as $\rho(t) = |\psi(t)|^2$
- Each particle follows individual trajectories, subjected to $V + V_Q$
quantal/Bohmian trajectories
- Make measurement: no need for collapse of wave-function



(C. Phillipides, C. Dewdney, B. J. Hiley, *Quantum Interference and the Quantum Potential*, *il nuovo cimento B*, **52**, no.1 (1979) 15-28)

- Stern-Gerlach, hydrogen atom, Aharonov-Bohm, harmonic oscillator, and all others \rightarrow correct predictions
- Only dynamical input \rightarrow Schrödinger equation
- Classical limit: trajectories \rightarrow with $\hbar \rightarrow 0$ (interference disappears)
- L. de Broglie (1927), D. Bohm (1952), superconductivity, superfluidity, Bose-Einstein condensates (Gross-Pitaevskii equation)

Valid picture!

Quantum Raychaudhuri equation

$V \rightarrow V + V_Q/m$ in Raychaudhuri equation

$$\frac{d\theta}{dt} = -\frac{1}{3} \theta^2 - \sigma_{ab}\sigma^{ab} - \nabla^2 V + \underbrace{\frac{\hbar^2}{2m^2} \nabla^2 \left(\frac{1}{\mathcal{R}} \nabla^2 \mathcal{R} \right)}_{\substack{\text{quantum correction} \\ \text{attractive or repulsive?} \\ \text{focusing or defocusing?}}}$$

$$\psi = \psi_0 e^{-r^2/L^2} \rightarrow \frac{\hbar^2}{2m^2} \nabla^2 \left(\frac{1}{\mathcal{R}} \nabla^2 \mathcal{R} \right) \sim +\frac{1}{L^4} \leftarrow \text{Repulsive at short distances}$$

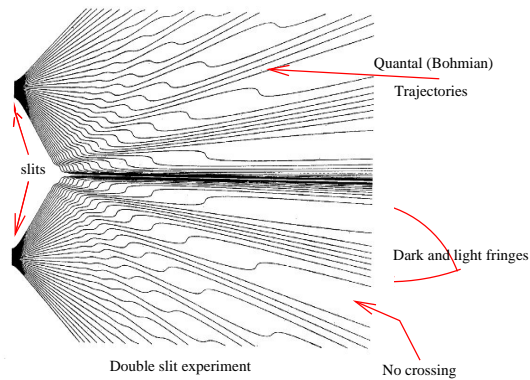
$$-\frac{\hbar^2}{2m} \nabla^2 \psi + [V + \underbrace{g|\psi|^2}_{\text{interaction}}] \psi = i\hbar \frac{\partial \psi}{\partial t} \quad (\text{Gross-Pitaevskii equation})$$

interaction

$$\psi = \psi_0 \tanh(r/L\sqrt{2}) \quad (g > 0), \quad \psi = \sqrt{2} \psi_0 \operatorname{sech}(r/L) \quad (g < 0)$$

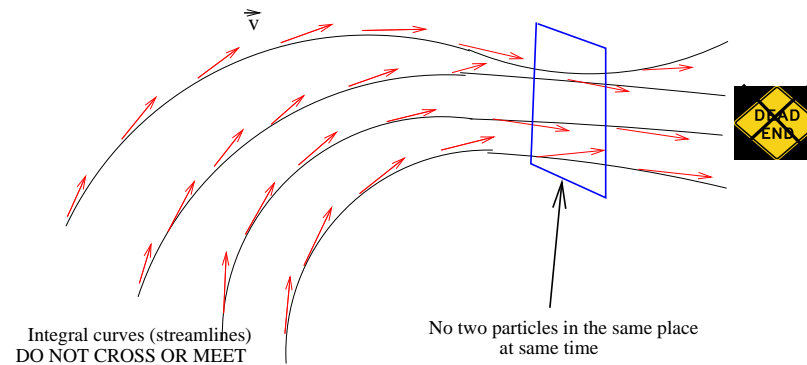
$$\frac{\hbar^2}{2m^2} \nabla^2 \left(\frac{1}{\mathcal{R}} \nabla^2 \mathcal{R} \right) \sim +\frac{1}{L^4} \leftarrow \text{Repulsive at short distances}$$

Repulsion may prevent focussing/conjugate points. But it gets better!



No-crossing of quantal (Bohmian) trajectories

$$\vec{v}(\vec{x}, t) = \frac{d\vec{x}}{dt} \equiv \frac{\hbar}{m} \mathcal{I}m \left(\frac{\vec{\nabla} \psi}{\psi} \right) = \frac{\hbar}{m} \vec{\nabla} S(\vec{x}, t) = \text{single valued}$$



Trajectories do not end in a caustic/focus. They go on forever.

Relativistic generalization

$$\left[\partial^2 + \frac{m^2 c^2}{\hbar^2} - \epsilon_1 R - \epsilon_2 \frac{i}{2} f_{cd} \sigma^{cd} \right] \Phi = 0$$

($\Phi = \mathcal{R} e^{iS}$, Normalizable, single valued)

$$k_a = \partial_a S, \quad u_a = c \frac{dx_a}{d\lambda} = \frac{\hbar k_a}{m}, \quad \vec{v} = \frac{d\vec{x}}{dt} = -c^2 \frac{\vec{\nabla} S}{\partial^0 S}$$

Imaginary part: $\partial^a (\mathcal{R}^2 \partial_a S) = \frac{\epsilon_2}{2} f_{cd} \sigma^{cd} \mathcal{R}^2$

Real part: $k^2 = \frac{(mc)^2}{\hbar^2} - \epsilon_1 R + \frac{\partial^2 \mathcal{R}}{\mathcal{R}}$

$$u^b{}_{;a} u^a = -\frac{\epsilon_1 \hbar^2}{m^2} R^{;b} + \frac{\hbar^2}{m^2} \left(\frac{\partial^2 \mathcal{R}}{\mathcal{R}} \right)^{;b}$$

(i.e. geodesic equation + $V_Q = \frac{\hbar^2}{m^2} \frac{\partial^2 \mathcal{R}}{\mathcal{R}}$)

Quantum Raychaudhuri equation

$$\frac{d\theta}{d\lambda} = -\frac{1}{3} \theta^2 - \sigma_{ab} \sigma^{ab} - R_{cd} u^c u^d$$

$$- \frac{\epsilon_1 \hbar^2}{m^2} h^{ab} R_{;a;b} - \frac{\hbar^2}{m^2} h^{ab} \left(\frac{\partial^2 \mathcal{R}}{\mathcal{R}} \right)^{;a;b}$$

↑ quantum corrections ↑

Null geodesics: $\frac{d\theta}{d\lambda} = -\frac{1}{3} \theta^2 - \sigma_{ab} \sigma^{ab} - R_{cd} u^c u^d - \epsilon_1 \hbar^2 h^{ab} R_{;a;b} - \hbar^2 h^{ab} \left(\frac{\partial^2 \mathcal{R}}{\mathcal{R}} \right)^{;a;b}$

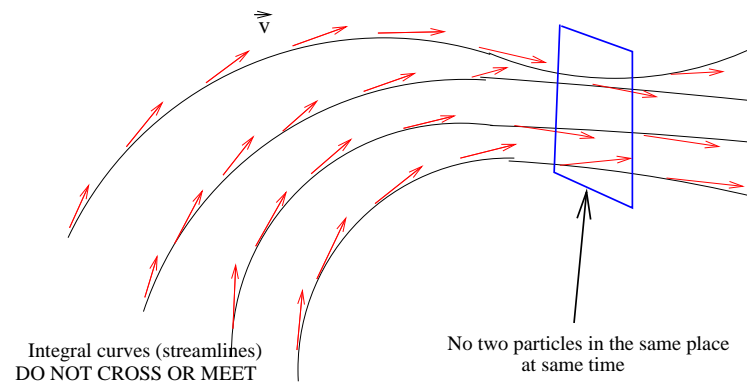
$$\left[\theta = Tr(k_{a;b}) \right]$$

Again

$$\Phi = \Phi_0 e^{-r^2/L^2}, \quad \Phi = \Phi_0 \tanh(r/L\sqrt{2}) \quad (g > 0), \quad \Phi = \sqrt{2} \Phi_0 \operatorname{sech}(r/L) \quad (g < 0)$$
$$\frac{\hbar^2}{m^2} \left(\frac{\partial^2 \mathcal{R}}{\partial \mathcal{R}} \right) \sim +\frac{1}{L^4} \quad \leftarrow \text{Repulsive at short distances. But, better...}$$

No-crossing of quantal (Bohmian) trajectories

$$\vec{v} = \frac{d\vec{x}}{dt} = -c^2 \frac{\vec{\nabla} S}{\partial^0 S}$$



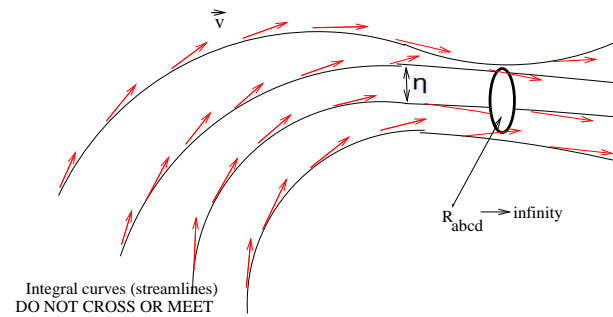
- No focusing, no conjugate points, geodesics go on forever
- No singularities! (all because of \hbar)

Curvature singularities

$$\frac{D^2 \eta^a}{d\lambda^2} = -R^a{}_{bfc} u^b u^c \eta^f - \frac{\hbar^2}{m^2} \left[\left(\frac{\partial^2 \mathcal{R}}{\mathcal{R}} \right)^{;a} \right]_{;c} \eta^c$$

↑ *quantum*

But $\vec{\eta} \neq 0$ anymore



Therefore, $R_{abcd} R^{abcd} \rightarrow \infty$ regions are not accessible!

Application

Cosmology (Raychaudhuri \rightarrow FRW)

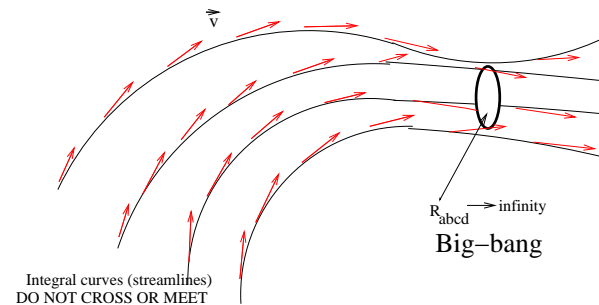
- Homogeneous, isotropic, $DE + DM + M$, GR cosmology, $a(t) =$ scale factor

$$\theta = 3 \frac{\dot{a}}{a} \quad , \quad R_{cd} u^c u^d = \frac{4\pi G}{3} (\rho + 3p) \quad \rightarrow$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3p) + \frac{\Lambda}{3} + \frac{\hbar^2}{3m^2} h^{ab} \left(\frac{\partial^2 \mathcal{R}}{\mathcal{R}} \right)_{;a;b}$$

(second order Friedmann equation + quantum correction)

- No crossing (e.g. at the big-bang)



- Addition of (quantum) cosmological constant/dark energy

$$\Lambda_Q = \frac{\hbar^2}{m^2 c^2} h^{ab} \left(\frac{\partial^2 \mathcal{R}}{\mathcal{R}} \right)_{;a;b}$$

$$\Lambda_Q$$

For wavefunctions considered earlier

$$(\psi = \psi_0 e^{-r^2/L^2}, \psi = \psi_0 \tanh(r/L\sqrt{2}) (g > 0), \psi = \sqrt{2} \psi_0 \operatorname{sech}(r/L) (g < 0))$$

$$\Lambda_Q = \frac{1}{L^2} = \frac{1}{(\text{Compton wavelength})^2} = \left(\frac{mc}{h}\right)^2$$

$$L = 1.4 \times 10^{26} \text{ m (Hubble radius/size of observable universe)}$$

$$m \approx 10^{-68} \text{ kg} = 10^{-32} \text{ eV}, \quad F = -\frac{Gm_1 m_2}{r^2} e^{-r/L}$$

$$\Lambda_Q = 10^{-52} \text{ m}^{-2} = 10^{-123} (\text{in Planck units}) = H_0^2/c^2$$

(Observed, also coincidence problem)

$m =$ graviton (or photon) mass?

Consistent with observations and theory

S. Das, arXiv:1405.4011

A. F. Ali, S. Das, arXiv:1404.3093

Summary

- Classical \rightarrow quantal (Bohmian) trajectories
 \Rightarrow Classical \rightarrow quantum Raychaudhuri equation
- No conjugate points, no geodesic incompleteness, no counterpart of singularity theorems for quantal trajectories
- Cosmology - addressing smallness and coincidence problems *and* no big bang
- Violation of the Equivalence principle? Quantum modifications of Einstein equations?
- Classical \rightarrow *quantum* spacetime: results still expected to hold